# Engineering Casimir forces with metamaterials

Diego A. R. Dalvit
Theoretical Division
Los Alamos National Laboratory







### Outline of this talk



- Brief review of theory and experiments on the Casimir force
- Materials effects: Casimir repulsion with metamaterials
- Conclusions

#### Collaborators

Theory: Peter Milonni (LANL)

Felipe da Rosa (LANL)

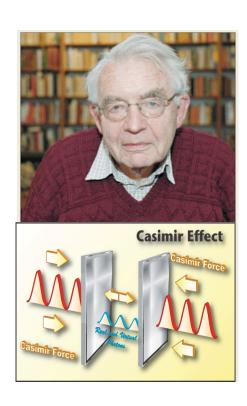
Experiment: Antoniette Taylor (CINT, LANL)

Steve Lamoreaux (Yale)

Ricardo Decca (Indiana)

### The Casimir force





Casimir forces originate from changes in quantum vacuum fluctuations imposed by surface boundaries

They were predicted by the Dutch physicist Hendrik Casimir in 1948

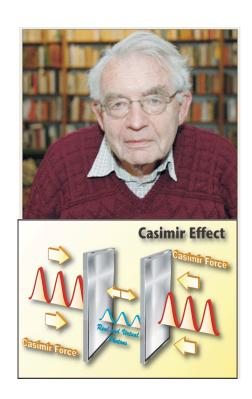
Dominant interaction in the micron and sub-micron lengthscales

$$\frac{F}{A} = \frac{\pi^2}{240} \; \frac{\hbar c}{d^4}$$

 $(130 \text{nN/cm}^2 @ d = 1 \mu\text{m})$ 

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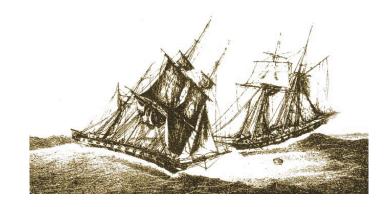
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Classical Analog: L'Album du Marin (1836)



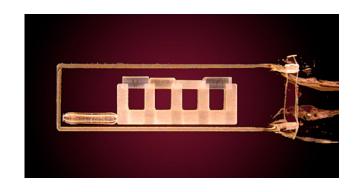
# Relevant applications



### Quantum Science and Technology:

Atom-surface interactions

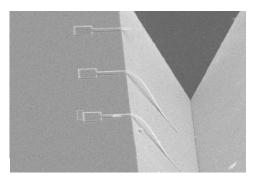
Precision measurements



Cornell et al (2007)

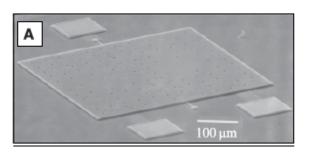
### Nanotechnology:

Problems with stiction of movable parts in MEMS



Zhao et al (2003)

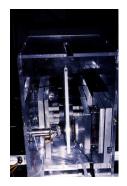
Actuation in NEMS and MEMS driven by Casimir forces

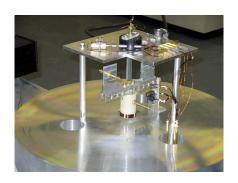


Capasso et al (2001)



### Torsion pendulum



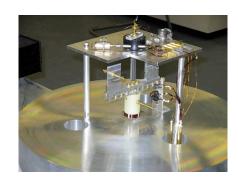


sphere-plane, d=1-10 um Lamoreaux



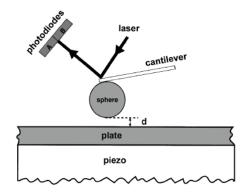
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### Atomic force microscope

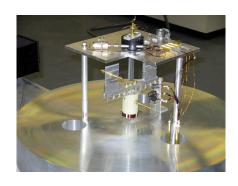


sphere-plane, d=200-1000 nm Mohideen et al



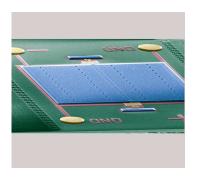
#### Torsion pendulum





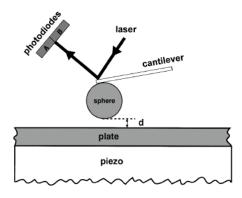
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#### MEMS and NEMS



sphere-plane, d=200-1000 nm Capasso et al, Decca et al

### Atomic force microscope



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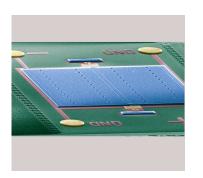
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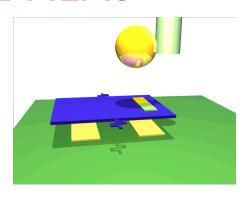




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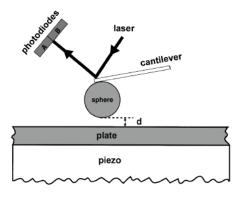
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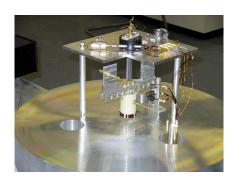


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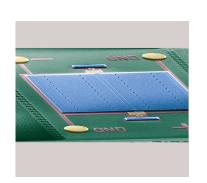
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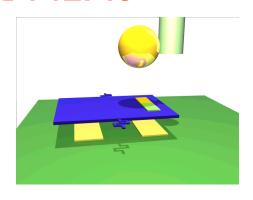




sphere-plane, d=1-10 um Lamoreaux

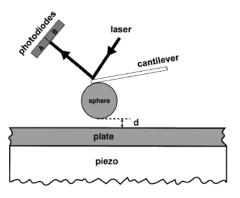
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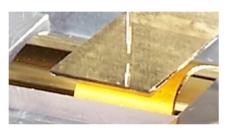
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### Atomic force microscope



sphere-plane, d=200-1000 nm Mohideen et al

#### Micro-cantilever



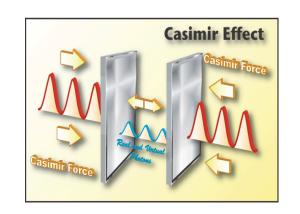
plane-plane, cylinder-plane, d=1-3 um Onofrio et al

# Tailoring the Casimir force



Magnitude and sign of the Casimir force depend on geometry and materials

Lifshitz formula: (assumes continuous and isotropic media)



# $\frac{F}{A} = 2k_B T \sum_{n=0}^{\infty} \int_{\xi_n/c}^{\infty} \frac{d\kappa}{2\pi} \kappa^2 \sum_{\substack{\text{TE-TM}}} \left( \frac{e^{2\kappa d}}{r_{\lambda_1} r_{\lambda_2}} - 1 \right)^{-1}$ $r_{\lambda}(i\xi_n)$ $\omega_n = i\xi_n = 2\pi i n$

Reflection coefficients:

$$r_{\lambda}(i\xi_n)$$

$$\omega_n = i\xi_n = 2\pi i n k_B T/\hbar$$

Reflection coefficients at imaginary frequencies --> Kramers-Kronig



$$\epsilon(i\xi_n) = 1 + \frac{2}{\pi} \int_0^\infty \frac{\omega \epsilon''(\omega)}{\omega^2 + \xi_n^2} d\omega \qquad \qquad \mu(i\xi_n) = 1 + \frac{2}{\pi} \int_0^\infty \frac{\omega \mu''(\omega)}{\omega^2 + \xi_n^2} d\omega$$

igspace The gap d sets a cut-off frequency:  $i\xi_{
m cut-off} \simeq c/d$ 

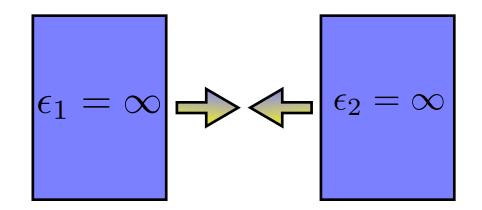
For  $d = 200 \text{nm} - 1 \mu \text{m}$  frequencies in the near-infrared/optical



#### Ideal attractive limit

Casimir 1948

$$\frac{F}{A} = +\frac{\pi^2}{240} \frac{\hbar c}{d^4}$$

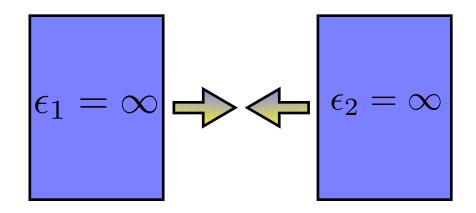




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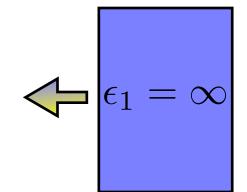
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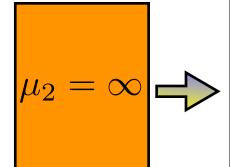


Ideal repulsive limit

Boyer 1974

$$\frac{F}{A} = -\frac{7}{8} \frac{\pi^2}{240} \frac{\hbar c}{d^4}$$



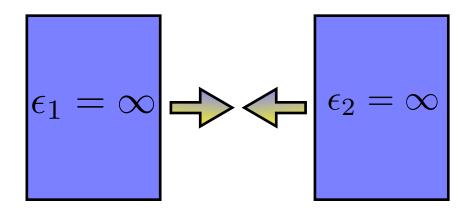




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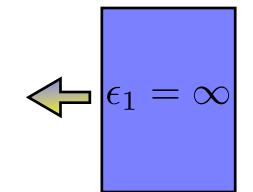
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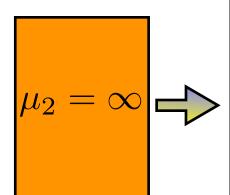


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Real repulsive limit  $\epsilon(i\xi) < \mu(i\xi)$ 

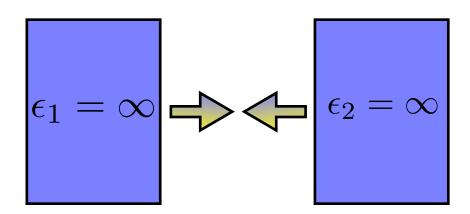
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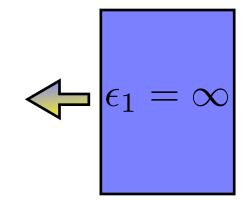
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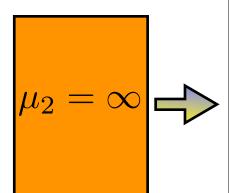


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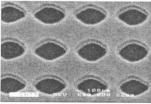


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Physicists have 'solved' mystery of levitation - Telegraph

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http://www.telegraph.co.uk/news/main.jhtml?xml=/news/2007/08/0...

Tuesday 4 September 2007

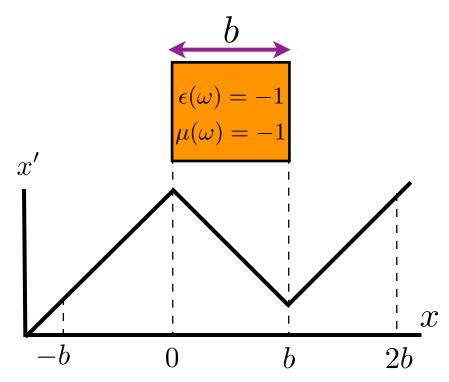


"In theory the discovery could be used to levitate a person"



#### Transformation media

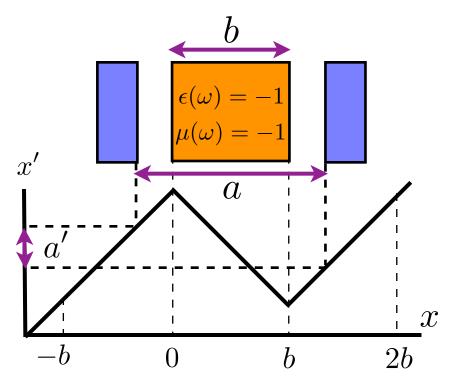
Leonhardt et al (2007)



Perfect lens: EM field in -b < x < 0 is mapped into x'. There are two images, one inside the device and one in b < x < 2b.



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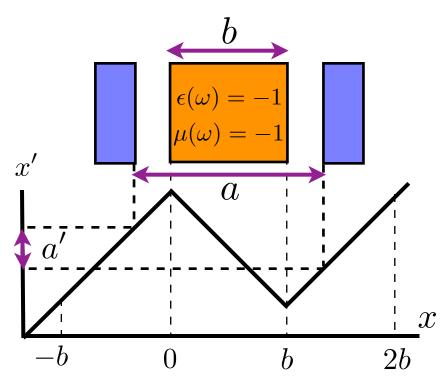
Casimir cavity: 
$$a' = |a - 2b|$$

When a < 2b (plates within the imaging range of the perfect lens)

$$\Rightarrow f = -\frac{\partial U}{\partial a'} \frac{\partial a'}{\partial a} = +\frac{\hbar c \pi^2}{240 a'^4} \Rightarrow \text{Repulsion}$$



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For real materials, however .....

- According to causality, no passive medium (  $\epsilon''(\omega) > 0$  ) can sustain  $\epsilon, \mu \simeq -1$ over a wide range of frequencies. In fact,  $\epsilon(i\xi), \mu(i\xi) > 0$
- Another proposal is to use an active MM (  $\epsilon$ " ( $\omega$ ) < 0) in order to get repulsion. But then the whole approach breaks down, as real photons would be emitted into the quantum vacuum.

### Metamaterials for Casimir



#### Drude-Lorentz model:

$$\epsilon_{\alpha}(\omega) = 1 - \frac{\Omega_{E,\alpha}^2}{\omega^2 - \omega_{E,\alpha}^2 + i\Gamma_{E,\alpha}\omega}$$
$$\mu_{\alpha}(\omega) = 1 - \frac{\Omega_{M,\alpha}^2}{\omega^2 - \omega_{M,\alpha}^2 + i\Gamma_{M,\alpha}\omega}$$

#### Typical separations

$$d = 200 - 1000 \text{ nm}$$

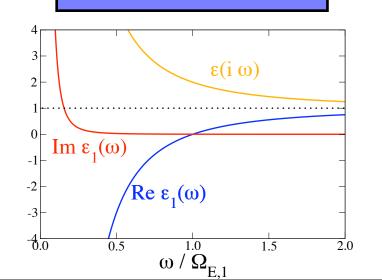


#### Infrared-optical frequencies

$$\Omega/2\pi = 5 \times 10^{14} \mathrm{Hz}$$

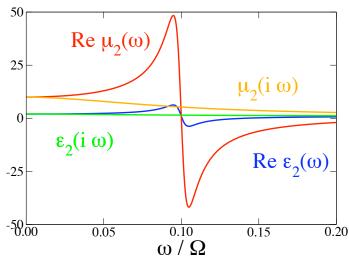
#### Drude metal (Au)

$$\Omega_E = 9.0 \; \mathrm{eV} \; \; \Gamma_E = 35 \; \mathrm{meV}$$



#### Metamaterial

Re 
$$\epsilon_2(\omega) < 0$$
 Re  $\mu_2(\omega) < 0$ 



$$\Omega_{E,2}/\Omega = 0.1 \quad \Omega_{M,2}/\Omega = 0.3$$

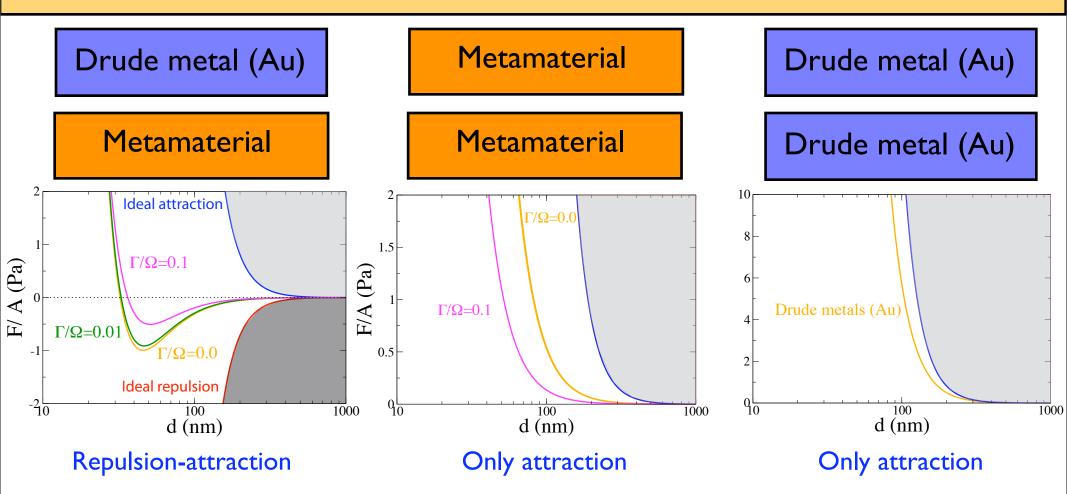
$$\omega_{E,2}/\Omega = \omega_{M,2}/\Omega = 0.1$$

$$\Gamma_{E,2}/\Omega = \Gamma_{M,2}/\Omega = 0.01$$

$$\epsilon(i\xi) < \mu(i\xi)$$

# Attraction-repulsion crossover





A slab made of Au ( $\rho=19.3~{\rm gr/cm^3}$ ) of width  $\delta=1\mu{\rm m}~{\rm could}$  levitate in front of one of these MMs at a distance of  $d\approx110~{\rm nm}~!!!$ 

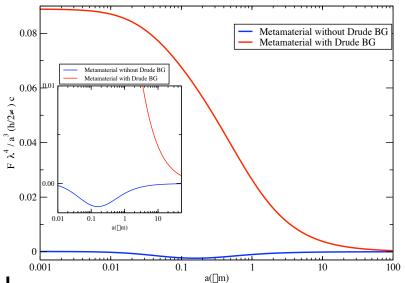
Casimir and metamaterials, Henkel et al (2005)
Casimir and surface plasmons, Intravaia et al (2005)
van der Waals in magneto-dielectrics, Spagnolo et al (2007)

# Some other important issues...



### Effects of Drude background

$$\epsilon(i\xi) = 1 + \underbrace{\frac{\Omega_D^2}{\xi^2 + \Gamma_D \xi}}_{\text{Drude part}} + \underbrace{\frac{\Omega_E^2}{\xi^2 + \omega_E^2 + \Gamma_E \xi}}_{\text{MM resonance}}$$



As the Drude background may overwhelm the resonant contribution in the low frequency limit, it may kill repulsion completely!



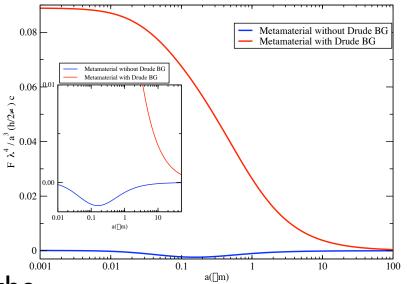
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$$\epsilon(i\xi) > \mu(i\xi)$$

### Effects of MM anisotropy

It is possible to derive a more complicated Lifshitz formula for continuous, anisotropic magneto-dielectric materials  $\epsilon_{ij}(\omega) = \mu_{ij}(\omega)$ 

$$\epsilon_{ij}(\omega) \quad \mu_{ij}(\omega)$$

Anisotropy typically reduces the magnitude of the possible Casimir repulsion, as compared to an ideally isotropic metamaterial

### Conclusions



- ☐ In principle, metamaterials can strongly influence the quantum vacuum, providing a route towards quantum levitation.
- ☐ However, we believe that previous works have been overly optimistic about the feasibility of quantum levitation via MMs.
- We have analyzed new important effects influencing Casimir repulsion in metamaterials:
  - Non-resonant optical response (Drude background)
  - Anisotropic permittivities and/or permeabilities
  - Different models for optical response (Drude, Drude-Lorentz, etc)